

# Left-right mixing effects on leptonic and semileptonic $b \rightarrow u$ decays

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- The dominant weak interaction in  $b \rightarrow u$  decays in the SM is strongly suppressed by the quark mixing matrix element  $|V_{ub}| \sim \lambda^4$  where  $\lambda \simeq 0.22$ .
- The relevant decays occur at the tree level, and are often sensitive to a non-standard physics beyond the SM if the new physics effects are not directly proportional to the small weak mixing.
- One of the simplest extensions of the SM corresponding to such a scenario is the general left-right model (LRM) with gauge group  $SU(2)_L \times SU(2)_R \times U(1)$ .  
(J. C. Pati and A. Salam, Phys. Rev. D10 275 (1974))
- Although the new physics effects in the LRM are followed by suppression factors such as the  $W_L - W_R$  mixing angle  $\xi$ , such suppression could be compensated by the right-handed quark mixing matrix  $V^R$  if  $V^R \neq V^L$  (nonmanifest LRM).

- From the global analysis of muon decay measurements, the lower bound on  $\xi$  can be obtained without imposing discrete left-right symmetry as follows:

$$\xi \leq \frac{g_R}{g_L} \frac{M_{W_L}^2}{M_{W_R}^2} < 0.034 \frac{g_L}{g_R}$$

(S.-h. Nam, Phys. Rev. D66 055008 (2002))

- $|V_{ub}^R|$  could be as large as  $\lambda$  for the following types of  $V^R$ :

$$\begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}, \quad \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}, \quad \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{pmatrix}$$

(T.G. Rizzo, Phys. Rev. D 58 114014 (1998))

- Although the mixing angle  $\xi$  is small, the combined parameter  $\xi V_{ub}^R / V_{ub}^L$  could significantly contribute to the value of  $|V_{ub}|$  extracted from the data in  $b \rightarrow u$  decays.

- Four-Fermion interaction for  $b \rightarrow c$  semileptonic decays:

$$\mathcal{H}_{\text{eff}} = 2\sqrt{2}G_F V_{qb}^L [(\bar{q}_L \gamma_\mu b_L) + \xi_q (\bar{q}_R \gamma_\mu b_R)] (\bar{\ell}_L \gamma_\mu \nu_L),$$

where  $\xi_q \equiv \xi(g_R V_{qb}^R)/(g_L V_{qb}^L)$  and  $q = u, c$ .

- Total decay rate of  $B \rightarrow l\nu X_c$ :

$$\Gamma(b \rightarrow cl\nu) \simeq \frac{G_F^2 m_b^5 |V_{cb}^L|^2}{192\pi^3} \left[ (1 + r_c^2 \xi_g^2) f(x) - r_c \xi_g h(x) \right]$$

where  $x \equiv m_c/m_b$ ,  $f(x) \sim 1 + O(x^2)$ , and  $h(x) \sim 4x + O(x^3)$ .

- Following approximate bound can be obtained by the comparison of  $|V_{cb}|_{\text{incl}}$  and  $|V_{cb}|_{\text{excl}}$ :

$$\xi_c \approx 0.14 \pm 0.18$$

(M.B. Voloshin, Mod. Phys. Lett. A **12**, 1823 (1997))

- In the LRM,

$$|V_{ub}|_{incl} \simeq |V_{ub}| + O(\xi_g^2) \sim O(\lambda^4) \quad (\lambda \simeq 0.22)$$

- 2008 HFAG averages of the selected theoretical methods for  $|V_{ub}|_{incl}$ :

$$|V_{ub}|_{incl} \times 10^3 = \begin{cases} 4.48 \pm 0.16 \begin{smallmatrix} +0.25 \\ -0.26 \end{smallmatrix} & \text{(DGE)} \\ 3.78 \pm 0.13 \pm 0.24 & \text{(ADFR)} \\ 4.09 \pm 0.20 & \text{(Our average)} \end{cases}$$

(C.-H. Chen and S.-h. Nam, Phys. Lett. B666, 462 (2008))

- Current HFAG averages of the selected theoretical methods for  $|V_{ub}|_{incl}$ :

$$|V_{ub}|_{incl} \times 10^3 = \begin{cases} 4.46 \pm 0.16 \begin{smallmatrix} +0.18 \\ -0.17 \end{smallmatrix} & \text{(DGE)} \\ 4.16 \pm 0.14 \begin{smallmatrix} +0.25 \\ -0.22 \end{smallmatrix} & \text{(ADFR)} \\ 4.34 \pm 0.18 & \text{(Our average)} \end{cases}$$

- Recently obtained value of  $|V_{ub}|_{excl}$  from  $B \rightarrow \pi e \nu$  is smaller than those of  $|V_{ub}|_{incl}$ :

$$|V_{ub}|_{excl}^{LCSR} = (3.5 \pm 0.4 \pm 0.2 \pm 0.1) \times 10^{-3}$$

- In the LRM, we can roughly estimate the mixing parameter  $\xi_u$  from the mismatch between the values of  $|V_{ub}|_{incl}$  and  $|V_{ub}|_{excl}$ :

$$\begin{aligned} |V_{ub}|_{excl}^{\pi e \nu} &= |V_{ub}^L| |1 + \xi_u| \simeq |V_{ub}|_{incl} |1 + \xi_u| \\ \Rightarrow \text{Re}(\xi_u) &= -(0.14 \pm 0.12) \quad (2008) \\ &= -(0.19 \pm 0.11) \quad (\text{current update}) \end{aligned}$$

$\Rightarrow$  Difference between the values of  $|V_{ub}|$  obtained from inclusive and exclusive modes gets bigger  $\rightarrow$  Hint for new physics?

- BELLE and BABAR collaborations had first found evidence for the purely leptonic  $B^- \rightarrow \tau^- \bar{\nu}_\tau$  decays in 2007 and 2008:

$$\begin{aligned}
 Br(B^- \rightarrow \tau^- \bar{\nu}_\tau) &= \begin{cases} (1.79^{+0.56}_{-0.49} +^{0.46}_{-0.51}) \times 10^{-4} & \text{(BELLE)} \\ (1.2 \pm 0.4 \pm 0.3 \pm 0.2) \times 10^{-4} & \text{(BABAR)} \end{cases} \\
 &= (1.41^{+0.43}_{+0.42}) \times 10^{-4} \quad \text{(2008 HFAG average)}
 \end{aligned}$$

- The BABAR result is an average of two results,  $(0.9 \pm 0.6 \pm 0.1) \times 10^{-4}$  and  $(1.8^{+0.9}_{-0.8} \pm 0.4 \pm 0.2) \times 10^{-4}$ , and the latter one is newer.
- Our earlier 2008 estimate of the branching fraction was:

$$\begin{aligned}
 Br(B^- \rightarrow \tau^- \bar{\nu}_\tau) &= Br(B^- \rightarrow \tau^- \bar{\nu}_\tau)^{SM} |1 - \xi_u|^2 \\
 &= (1.78 \pm 0.53) \times 10^{-4}
 \end{aligned}$$

$\implies$  agreed very well with the BELLE result and the newer BABAR result, but not with the older BABAR result.

- Current HFAG average of new BELLE and BABAR measurements for the purely leptonic  $B^- \rightarrow \tau^- \bar{\nu}_\tau$  decays:

$$\begin{aligned} Br(B^- \rightarrow \tau^- \bar{\nu}_\tau) &= \begin{cases} (1.54^{+0.38}_{-0.37} + {}^{0.29}_{-0.31}) \times 10^{-4} & \text{(BELLE)} \\ (1.76 \pm 0.49) \times 10^{-4} & \text{(BABAR)} \end{cases} \\ &= (1.64 \pm 0.34) \times 10^{-4} \quad \text{(current HFAG average)} \end{aligned}$$

- Our current estimate of the branching fraction is:

$$\begin{aligned} Br(B^- \rightarrow \tau^- \bar{\nu}_\tau) &= Br(B^- \rightarrow \tau^- \bar{\nu}_\tau)^{SM} |1 - \xi_u|^2 \\ &= (1.72 \pm 0.37) \times 10^{-4} \end{aligned}$$

$\implies$  agrees very well with the current HFAG average.



- Differential decay rate for  $B \rightarrow \rho \ell \nu$  mode:

$$\frac{d^2\Gamma(B^0 \rightarrow \rho^- \ell^+ \nu_\ell)}{dy d\cos\theta_\ell} = \frac{G_F^2 m_B^2 |\mathbf{p}_\rho| y}{256\pi^3} |V_{ub}^L|^2 [(1 - \cos\theta_\ell)^2 |H_+|^2 + (1 + \cos\theta_\ell)^2 |H_-|^2 + 2\sin\theta_\ell^2 |H_0|^2],$$

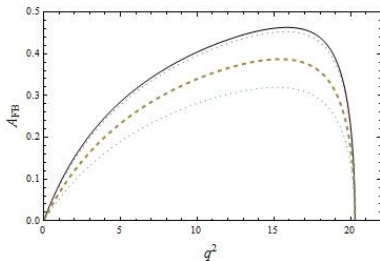
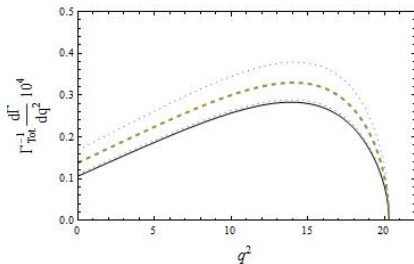
where the three helicity amplitudes  $H_{\pm,0}$  depend effectively on only three form factors  $V$  and  $A_{1,2}$ :

$$H_{\pm} = \frac{1}{m_B + m_\rho} \left[ (m_B + m_\rho)^2 (1 - \xi_u) A_1(q^2) \mp 2m_B |\mathbf{p}_\rho| (1 + \xi_u) V(q^2) \right],$$

$$H_0 = \frac{m_B(1 - \xi_u)}{2m_\rho(m_B + m_\rho)\sqrt{y}} \left[ \left( 1 - \frac{m_\rho^2}{m_B^2} - y \right) (m_B + m_\rho)^2 A_1(q^2) - 4|\mathbf{p}_\rho|^2 A_2(q^2) \right]$$

## Other semi-leptonic $b \rightarrow u$ transitions

- $d\Gamma(B^0 \rightarrow \rho^- \ell^+ \nu_\ell)/dq^2$  distribution and  $A_{\text{FB}}(B^0 \rightarrow \rho^- \ell^+ \nu_\ell)$  as a function of  $q^2$  for the SM (solid line),  $\xi_u^r = -0.14$  (dashed line), and its error (dotted line).



- Differential decay rate for  $B \rightarrow \gamma \ell \nu_\ell$  mode:

$$\frac{d^2\Gamma(B^- \rightarrow \gamma \ell^- \bar{\nu}_\ell)}{dy d\cos\theta} = \frac{\alpha_{em} G_F^2 m_B^5}{512\pi^2} y(1-y)^3 |V_{ub}^L|^2 (1 - \hat{m}_\ell^2)^2 I(q^2, \cos\theta),$$

with

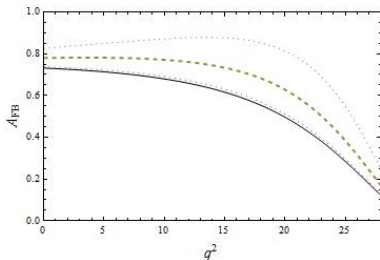
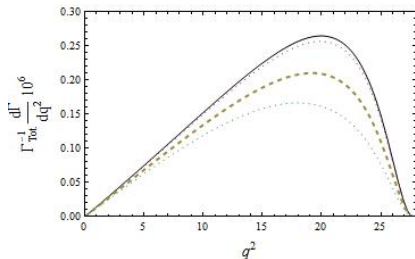
$$\begin{aligned} I(q^2, \cos\theta) &= |F'_V + F'_A|^2 \left[ 1 + \hat{m}_\ell^2 + (1 - \hat{m}_\ell^2) \cos\theta \right] (1 + \cos\theta) \\ &\quad + |F'_A - F'_V|^2 \left[ 1 + \hat{m}_\ell^2 - (1 - \hat{m}_\ell^2) \cos\theta \right] (1 - \cos\theta), \end{aligned}$$

where

$$F'_V = F_V(1 + \xi_u) \quad F'_A = F_A(1 - \xi_u).$$

## Other semi-leptonic $b \rightarrow u$ transitions

- $d\Gamma(B^- \rightarrow \gamma \ell^- \bar{\nu}_\ell)/dq^2$  distribution and  $A_{\text{FB}}(B^- \rightarrow \gamma \ell^- \bar{\nu}_\ell)$  as a function of  $q^2$  for the SM (solid line),  $\xi_u^r = -0.14$  (dashed line), and its error (dotted line).



# Summary

- Admixture of a right-handed  $b \rightarrow u$  current could give a significantly different contributions to the inclusive and exclusive rates of the semileptonic decays of the  $B$  mesons.
- From the current mismatch between  $|V_{ub}|_{incl}$  and  $|V_{ub}|_{excl}$  obtained from the independent experiments, we estimate the size of the left-right mixing parameter  $\xi_u$  to be  
 $\text{Re}(\xi_u) = -(0.14 \pm 0.12)$  (2008 prediction, published)  
 $\text{Re}(\xi_u) = -(0.19 \pm 0.11)$  (current update).
- For  $\text{Re}(\xi_u) = -0.14$ , we show that the branching fraction for leptonic  $B \rightarrow \tau\nu$  and semileptonic  $B \rightarrow \rho\ell\nu$  decays can be enhanced by 30% and 17%, respectively, while the branching fraction for radiative leptonic  $B \rightarrow \gamma\ell\nu$  decays can be reduced by 18%.
- Since the left-right mixing contributions obtained in this letter in leptonic and semileptonic  $b \rightarrow u$  decays are not simply negligible, our estimate could be a reasonable guide to search for the existence of the right-handed current.